

# Extensions of brackets in classical field theories

Summer School of Geometry, Dynamics and Field Theory 2025

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ICMAT-UNIR

# Structure of the talk

Introduction to the problem

Geometric stage and algebraic structure of observables

Hamiltonians and extensions of brackets

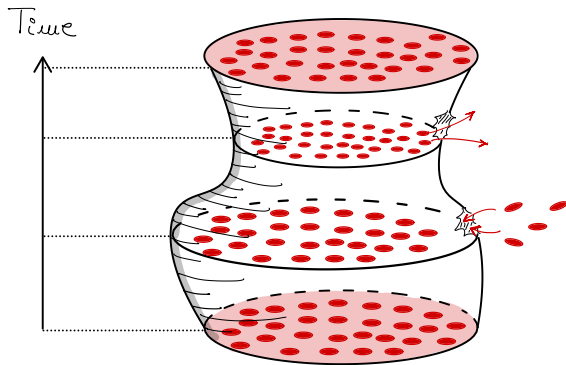
References

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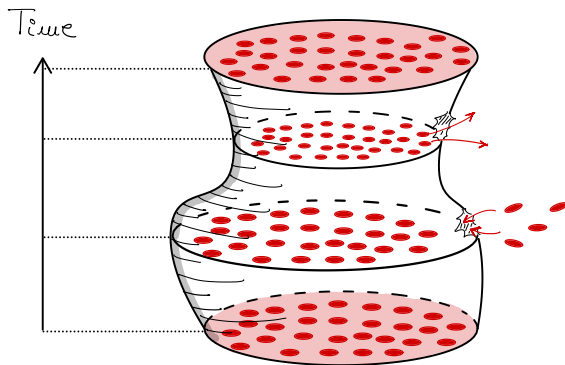
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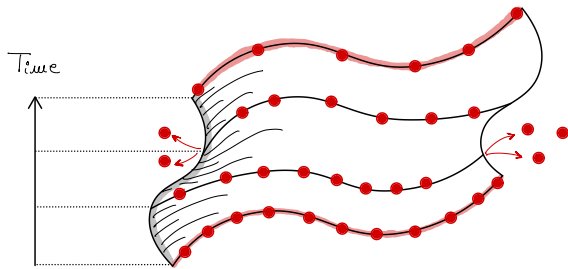


For  $\alpha \in \Omega^{n-1}(M)$ :

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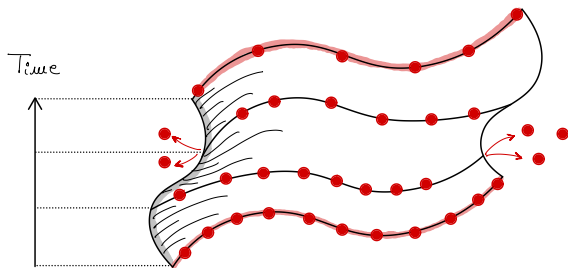
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Problem (that we solve): Determine evolution of forms via some bracket:

$$\psi^*(d\alpha) = d\alpha + \{\alpha, \mathcal{H}\},$$

where  $\psi$  is a solution of the corresponding PDE and  $\alpha$  is some  $a$ -form.

## Previous work:

1. Igor V. Kanatchikov. **“Canonical Structure of Classical Field Theory in the Polymomentum Phase Space”**. In: *Rep. Math. Phys.* **41.1** (1998), pp. 49–90
2. Miguel Á. Berbel and Marco Castrillón-López. **“Poisson–Poincaré Reduction for Field Theories”**. In: *J. Geom. Phys.* **191** (2023), p. 104879
3. François Gay-Balmaz, Juan C. Marrero, and Nicolás Martínez-Alba. **“A New Canonical Affine Bracket Formulation of Hamiltonian Classical Field Theories of First Order”**. In: *Rev. Real Acad. Cienc. Exactas Fis. Nat. Ser. A-Mat.* **118.3** (2024), p. 103

# Geometric stage and algebraic structure of observables

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# General setup for field equations

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- (i)  $\alpha_1, \dots, \alpha_k \in \Omega^{n-1}(M)$  be **semi-basic forms** (representing observables).
- (ii)  $\beta_1, \dots, \beta_k \in \Omega^n(M)$  be **basic forms** (representing evolution of observables).



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We deal with partial differential equations with the following structure\*:

$$\psi^*(d\alpha_i) = \beta_i \circ \psi, \text{ where } \psi : X \rightarrow M \text{ is a section.}$$

## Examples (equations)

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(i) **Hamilton equations:**

(a) Fiber bundle:  $T^*Q \times \mathbb{R} \rightarrow \mathbb{R}$ .

(b) Semi-basic forms:  $q^i, p_i, t$ .

(c) Basic forms:  $\frac{\partial H}{\partial p_i} dt, -\frac{\partial H}{\partial q^i} dt, dt$ .

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  - (a) Fiber bundle: (covariant phase space)  $\wedge_2^n Y / \wedge_1^n Y \rightarrow X$ .
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(a) Fiber bundle:  $\text{Im } \text{leg}_{\mathcal{L}} \rightarrow X$ .

(b) Semi-basic forms:

$$A_\mu^i d^{n-1}x_\nu - A_\nu^i d^{n-1}x_\mu, F_i^{\mu\nu} d^{n-1}x_\nu, \frac{1}{n} x^\mu d^{n-1}x_\mu.$$

(c) Basic forms:

$$\left( F_{\mu\nu}^i - f_{jk}^i A_\nu^j A_\mu^k \right) d^n x, \left( -f_{jk}^i F_i^{\mu\nu} A_\mu^k \right) d^n x, d^n x.$$

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The observables  $\alpha_1, \dots, \alpha_k \in \Omega^{n-1}(M)$  allow us to define the space of **Hamiltonian forms**:

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**Assumption 1:** There is a bracket  $\{\cdot, \cdot\}$  on the space of Hamiltonian forms satisfying the following properties:

- (i) It is *skew-symmetric*:  $\{\alpha, \beta\} = -\{\beta, \alpha\}$ .
- (ii) It satisfies the Jacobi identity up to an exact term:

$$\{\alpha, \{\beta, \gamma\}\} + \{\beta, \{\gamma, \alpha\}\} + \{\gamma, \{\alpha, \beta\}\} = \text{exact form}.$$

- (iii) It vanishes on closed forms:  $d\alpha = 0 \implies \{\alpha, \beta\} = 0$ .
- (iv) There is a correspondence  $\alpha \mapsto X_\alpha$  such that
$$\{\alpha, \beta\} = \iota_{X_\beta} d\alpha.$$



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$$\{A_\mu^i d^{n-1}x_\nu - A_\nu^i d^{n-1}x_\mu, F_j^{\alpha\beta} d^{n-1}x_\beta\} = \delta_j^i \delta_{\mu\nu}^{\alpha\beta} d^{n-1}x_\beta.$$

# Hamiltonian forms of arbitrary order I

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## Definition

We say that a form  $\alpha \in \Omega^a(M)$  is **special Hamiltonian** if there is a semi-basic form  $\beta \in \Omega^{a+1}(M)$  such that  $\psi^*(d\alpha) = \beta \circ \psi$ , for every solution of the equations. The space of special Hamiltonian  $a$ -forms is denoted by  $\tilde{\Omega}_H^a(M)$ .

## Remark

*If  $\alpha \in \Omega^{n-1}(M)$  is special Hamiltonian,  $\alpha \in \Omega_H^{n-1}(M)$ .*

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## Theorem (de León, I.L. 2025)

*$\alpha \in \Omega^a(M)$  is special Hamiltonian if and only if  $\alpha \wedge \varepsilon \in \Omega_H^{n-1}(M)$ , for every closed and basic  $\varepsilon \in \Omega^{n-1-a}(M)$ .*

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A form  $\alpha \in \Omega_H^a(M)$  is called **Hamiltonian** if

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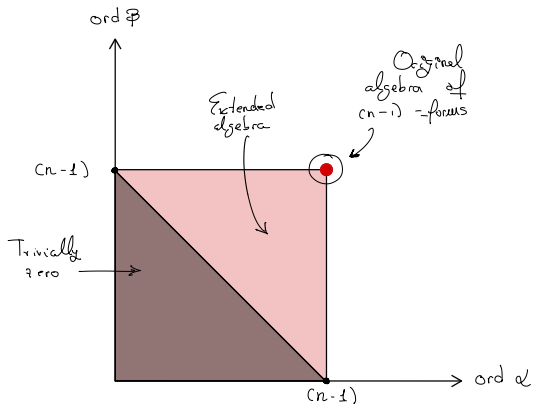
## Theorem (de León, I.L. 2025)

*There is an unique induced graded Poisson bracket*

$$\Omega_H^a(M) \otimes \Omega_H^b(M) \rightarrow \Omega_H^{a+b-(n-1)}(M)$$

*that maintains the properties of the original bracket of  $(n-1)$ -forms. Furthermore\*, special Hamiltonian forms define a subalgebra.*

# Summary



# Hamiltonians and extensions of brackets

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# Domain of definition of current brackets

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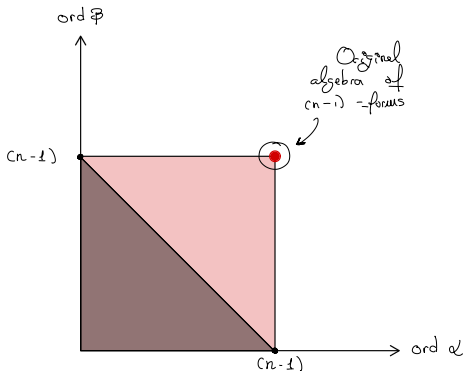
$$\begin{cases} \psi^*(d\alpha) = d\alpha + \{\alpha, \mathcal{H}\} \\ \deg\{\alpha, \mathcal{H}\} = \deg \alpha + \deg \mathcal{H} - (n - 1) \end{cases} \implies \deg \mathcal{H} = n.$$

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# First extension of the brackets I

## Theorem (de León, I.L. 2025)

*There exists an unique extension of  $\{\cdot, \cdot\}$*

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*for arbitrary  $a \geq 0$  that satisfies the properties of  $\{\cdot, \cdot\}$ .*



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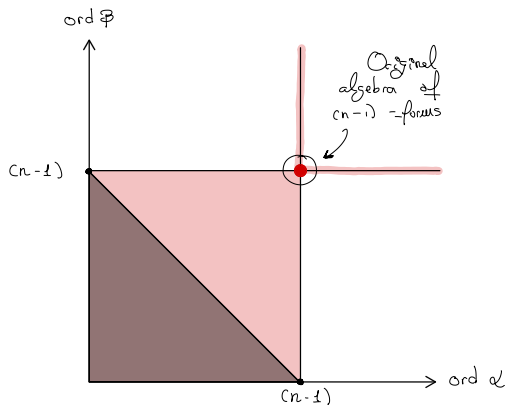
**Assumption 2:** There exists a form  $\mathcal{H} \in \Omega^n(M)[1]$ , the Hamiltonian, such that

$$\psi^*(d\alpha) = d\alpha + \{\alpha, \mathcal{H}\},$$

for every solution  $\psi$  and  $\alpha \in \Omega_H^{n-1}(M)$ .

# First extension of the brackets II

Current domain of definition:



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$$\mathcal{H} = \left( -\frac{1}{4} F_i^{\mu\nu} F_{\mu\nu}^i + \frac{1}{2} f_{jk}^i F_i^{\mu\nu} A_\mu^j A_\nu^k \right) d^n x - F_i^{\mu\nu} dA_\mu^i \wedge d^{n-1} x_\nu.$$

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Hamiltonian = Poincaré–Cartan form

## Final extension of the bracket I

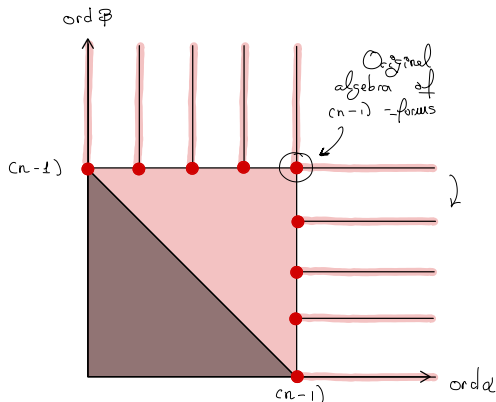


# Final extension of the bracket I

**Question:** Can we interpret  $\psi^*(d\alpha) = d\alpha + \{\alpha, \mathcal{H}\}$ , for arbitrary  $\alpha \in \Omega_H^a(M)$ ?

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## Final extensions of the brackets II

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**Theorem (de León, I.L. 2025)**

*There is a bijective correspondence between the possible extensions of the bracket and affine maps*

$$\gamma : \{\text{Hamiltonians}\} \rightarrow \{\text{Ehresmann connections on } \tau : M \rightarrow X\}$$

*such that  $\gamma(\mathcal{H})$  solves the Hamilton–De Donder–Weyl equations of  $\mathcal{H}$ , for every  $\mathcal{H}$ .*

# Final extension of the bracket III

## Corollary

*Let  $\gamma$  be such a map,  $\mathcal{H}$  be a Hamiltonian, and  $\psi$  be an integral section of  $\gamma(\mathcal{H})$ . Then*

$$\psi^*(d\alpha) = d\alpha + \{\alpha, \mathcal{H}\}_\gamma,$$

*for every  $\alpha \in \Omega_H^a(M)$ .*

# Final extension of the bracket III

## Corollary

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## Corollary

*Let  $\alpha$  be a special Hamiltonian form. Then, the bracket  $\{\alpha, \mathcal{H}\}_\gamma$  is independent of extension  $\{\cdot, \cdot\}_\gamma$ , for every Hamiltonian  $\mathcal{H}$ .*

## Technical remarks

The construction was based on a generalization of the  $\sharp$  mapping associated to a graded Poisson bracket. In particular, we generalized the techniques employed in

1. Peter W. Michor. “**A Generalization of Hamiltonian Mechanics**”. In: *J. Geom. Phys.* **2.2** (1985), pp. 67–82
2. Janusz Grabowski. “**Z-Graded Extensions of Poisson Brackets**”. In: *Rev. Math. Phys.* **09.01** (1997), pp. 1–27

to extend the brackets.

## Final remarks and remaining questions

- (i) The previous theoretical results seem to indicate that the subalgebra of special Hamiltonian forms is of high relevance to a particular field theory. We would like to compute these subalgebras for several almost regular Lagrangians to further study these classical field theories.
- (ii) We would also like to investigate the relation between these extensions and the instantaneous split formalism.
- (iii) It is also interesting to investigate the implications of this algebraic structure in the study of momentum maps and reduction, employing the graded brackets.



# References

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# Main references

1. Manuel de León and Rubén Izquierdo-López. **“Graded Poisson and Graded Dirac Structures”**. In: *J. Math. Phys.* **66.2** (2025). 10.1063/5.0243128, p. 022901
2. de León, Manuel and Izquierdo-López, Rubén. ***A description of classical field equations using extensions of graded Poisson brackets***. To appear soon in arXiv. 2025

Thank you for your attention!